Diagrammatic Monte Carlo approach to angular momentum in quantum many-body systems

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IMPURITIES AND ANGULAR MOMENTUM

Motivation for the study of composite impurities, i.e. impurities possessing angular momentum, comes from many different fields. They can be realized as:

- Molecules embedded into helium nanodroplets.
- Ultracold molecules and ions.
- Electronic excitations in Rydberg
 atoms.
- Angular momentum transfer from **electrons** to a **crystal lattice**.



DIAGRAMS AND UPDATES

Moving particle: *linear* momentum circulating on lines.



Feynman rules

Each free propagator $\lambda_i \mu_i$

Rotating particle: angular momentum circulating on lines.



Image from: J. P. Toennies and A. F. Vilesov, Angew. Chem. Int. Ed. 43, 2622 (2004).

THE ANGULON



The **angulon quasiparticle**: a quantum rotor dressed by a field of many-body excitations.

Each phonon propagator $\sum_{\lambda_i \mu_i} (-1)^{\mu_i} D_{\lambda_i}$ $\lambda_i \mu_i$ Each vertex $(-1)^{\lambda_i} \langle \lambda_i | |Y^{(\lambda_j)}| | \lambda_k \rangle \begin{pmatrix} \lambda_i & \lambda_j & \lambda_k \\ \mu_i & \mu_j & \mu_k \end{pmatrix}$

 $\sum_{\lambda_i \mu_i} (-1)^{\mu_i} G_{0,\lambda_i}$

The configuration space is **larger** than that of the Fröhlich polaron: in the diagram above j and λ can couple to give j' in many different ways. The configuration space is also **essentially different**: consider the second diagram below, angular momentum is not conserved on each phonon line (i.e. a phonon line subtracts 0 quanta of angular momentum, but gives back 2)...





....but it must be conserved on each 1-particle-irreducible component.



Solution

DIAGRAMMATIC MONTE CARLO

General structure of an impurity problem: $\hat{H} = \hat{H}_{imp} + \hat{H}_{bos} + \hat{H}_{imp-bos}$



DiagMC idea: set up a stochastic process sampling among all diagrams. **Configuration space:** diagram topology, phonons internal variables, times... The number of variables varies with the topology!

How: ergodicity, detailed balance $w_1p(1 \rightarrow 2) = w_2p(2 \rightarrow 1)$, achieved with a two-step update (proposal+acceptance/rejection).

Result: each configuration is visited with probability \propto its weight.

Works in continuous time and in the thermodynamic limit: no finite-size

Shuffle update: select one 1-particle-irreducible component, then shuffle the momenta couplings to another allowed configuration, keeping phonon momenta fixed but including configurations in which a phonon arc sub-tracts/gives back a different number of quanta.

The Add and Remove updates now only act on phonon lines which add/subtract the same number of quanta of angular momentum.

This scheme allows us to **visit all diagrams**.

Results

The ground-state energy, the energy of the first two excited states and the quasiparticle weight for the angulon are obtained by fitting the long-imaginary-time behaviour of G_j with

$$G_j(\tau) = Z_j \exp(-E_j \tau)$$

as a function of the dimensionless bath density \tilde{n} . They are compared with the weak- and strong-coupling theories.



effects or systematic errors.

For the Fröhlich polaron three updates are enough to explore all possible diagrams:
✓Add update: adds a new arc to a diagram.
✓Remove update: removes an arc from the diagram.
✓Change update: modifies the total length of the diagram.

REFERENCES

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CONCLUSIONS

- A technique for **molecular simulations** using the DiagMC framework.
- Angular momentum and rotations can be described with DiagMC. The configuration space is bigger, and an additional update is needed.
- Works naturally in the **thermodynamic limit** and in **continuous time**: no finite-size effect, no systematic errors.
- Straightforward access to the **Green's function** and to angular momentum properties: they are encoded in the formalism.